

# THERMODYNAMICS

MIKLÓS KELLERMAYER

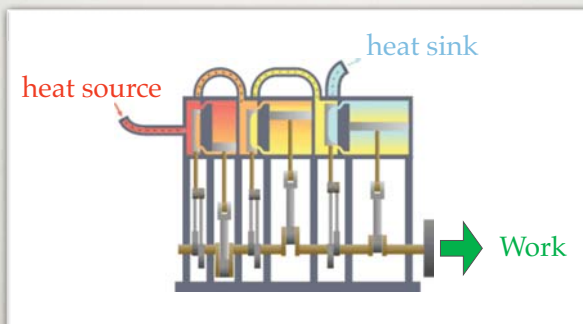
## OUTLINE

- Thermodynamic system.
- Characterization of the system. Thermodynamic variables.
- Microscopic structure of the system - internal energy.
- Equilibrium, reversibility.
- Zeroth Law of thermodynamics.
- Change, reversible change, thermodynamic processes.
- First Law of thermodynamics. Change in internal energy. Energy is conserved!
- Enthalpy.
- Second Law of thermodynamics.
- Direction of spontaneous processes. Entropy. Statistical interpretation. Thermodynamic probability. Phenomenological interpretation of entropy.
- Third Law of thermodynamics.
- Helmholtz- and Gibbs-free energy. Significance of free energy and entropy.
- Chemical potential.
- Irreversible thermodynamics, transport phenomena.
- Diffusion, Brown-movement.
- Osmosis.

## WHAT IS THERMODYNAMICS?

*Definition:* branch of physics studying the effects of changes in temperature, pressure and volume (energy changes).

*Therme* (Gr) heat, *dynamis* (Gr) power.



## THERMODYNAMICS

"Classical thermodynamics... is the only physical theory of universal content concerning which I am convinced that, within the framework of applicability of its basic contents, will never be overthrown."

-Albert Einstein

- Its importance stems from its generality.
- It reveals the driving forces behind natural phenomena.
- Due to its generalizations, it has abstract, complex and difficult-to-understand concepts.

"Thermodynamics is a funny subject. The first time you go through it, you don't understand it at all. The second time you go through it, you think you understand it, except for one or two small points. The third time you go through it, you know you don't understand it, but by that time you are so used to it, it doesn't bother you any more."

-Arnold Sommerfeld

# SHORT HISTORY OF THERMODYNAMICS



Otto von Guericke (1650) first vacuum pump.



Robert Boyle (1660) Boyle's law:  $pV = \text{constant}$ .



Joseph Louis Gay-Lussac (1802) universal gas law:  $pV = nRT$ .



Sadi Carnot (1824) father of thermodynamics. Concept of "work".



William Thomson (Lord Kelvin) (1849) name of "thermodynamics".



Rudolf Clausius (1850) concepts of enthalpy, entropy.



James Clerk Maxwell (1871) statistical thermodynamics, concept of equilibrium.



Ludwig Boltzmann (1875) statistical interpretation of entropy ( $S = k \ln W$ ).

Willard Gibbs (1876) concept of free energy (useful work).

Lars Onsager (1929) irreversible thermodynamics, Onsager's linear relationship.



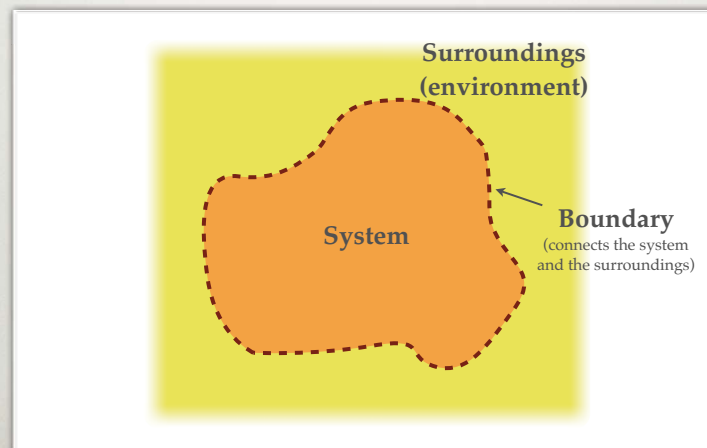
Christopher Jarzynski (1997) Jarzynski's equality. Free-energy change can be calculated from average irreversible work:  $e^{-\Delta G/k_B T} = \langle e^{-W/k_B T} \rangle$

# THERMODYNAMIC SYSTEM



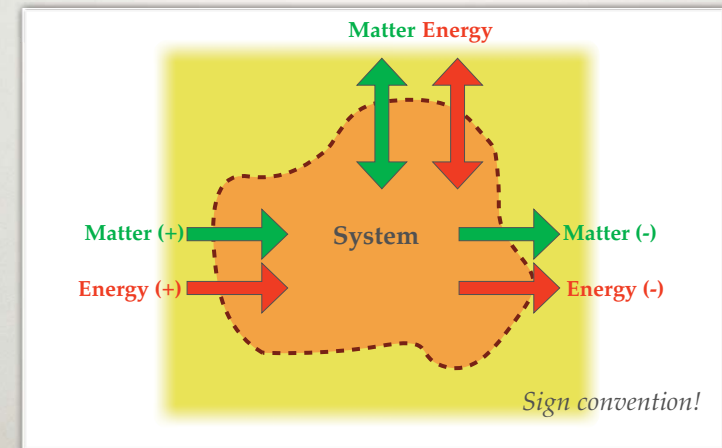
# ABSTRACTION OF THE THERMODYNAMIC SYSTEM

*Definition:* the thermodynamic system is the part of nature under investigation.



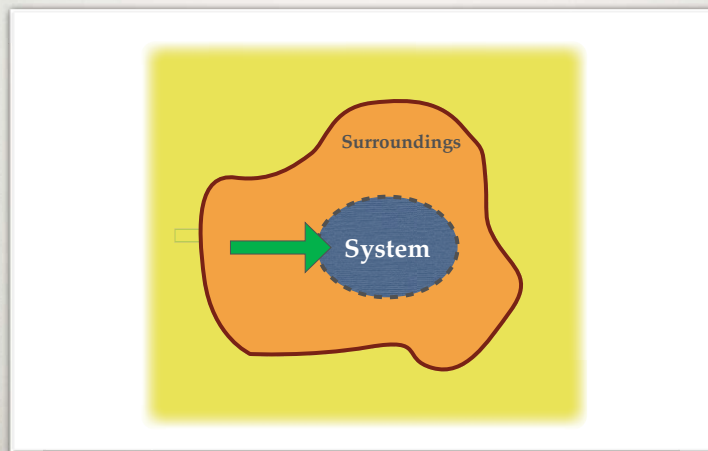
# THE THERMODYNAMIC SYSTEM INTERACTS WITH THE SURROUNDINGS

Exchange of matter and energy may occur across the boundary.

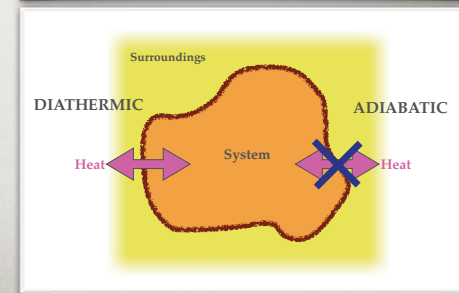
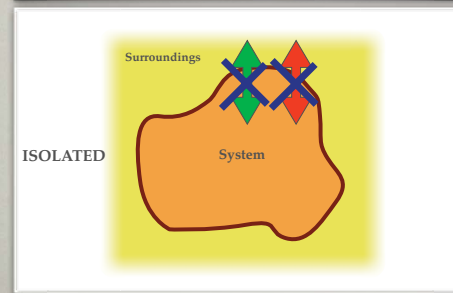
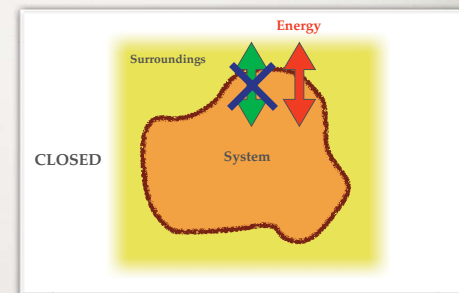
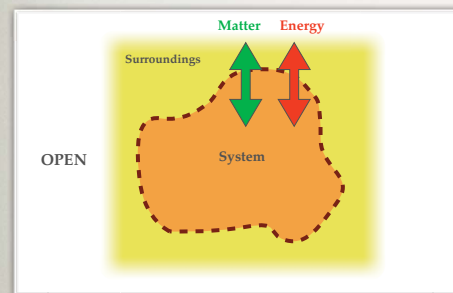


# THE THERMODYNAMIC SYSTEM INTERACTS WITH THE SURROUNDINGS

Defining the thermodynamic system has consequences for processes.



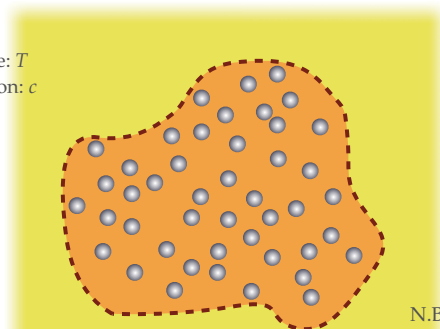
# TYPES OF THERMODYNAMIC SYSTEMS



# CHARACTERIZATION OF THE THERMODYNAMIC SYSTEM

Macroscopic characterization: state variables - explicitly determine the *state* of the system.

Pressure:  $p$   
Volume:  $V$   
Temperature:  $T$   
Concentration:  $c$



N.B.: universal gas law  
 $pV=nRT$

# CHARACTERIZATION OF THE THERMODYNAMIC SYSTEM

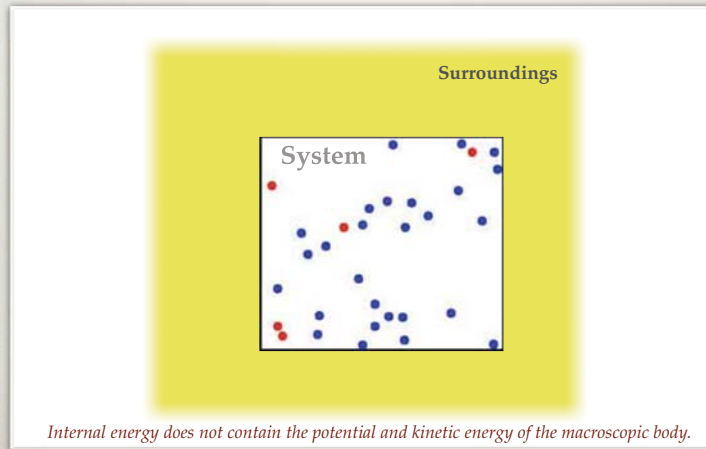
Macroscopic characterization: extensive and intensive variables  
Extensive variables: their value is proportional to the size of the system  
Intensive variables: their value is independent of the size of the system

Extensive variable	Intensive variable	Product: energy change
Volume ( $V$ ) Matter ( $n$ ) Charge ( $Q$ ) Entropy (disorder, $S$ )	Pressure ( $p$ ) Chemical potential ( $\mu$ ) Electric potential ( $\phi$ ) Temperature ( $T$ )	Volumetric work ( $p\Delta V$ ) Work of material transport ( $\mu\Delta n$ ) Electric work ( $\phi\Delta Q$ ) Heat ( $T\Delta S$ )

# CHARACTERIZATION OF THE THERMODYNAMIC SYSTEM

Microscopic characterization: internal energy ( $E$ )

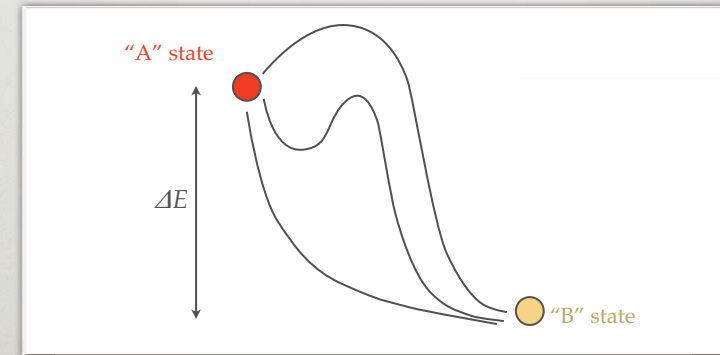
$$E = E_{el} + E_{vib} + E_{rot} + E_{other}$$



# INTERNAL ENERGY ( $E$ ) IS "STATE FUNCTION"

*State functions:* single-value functions of the state variables (independent variables) of the system.

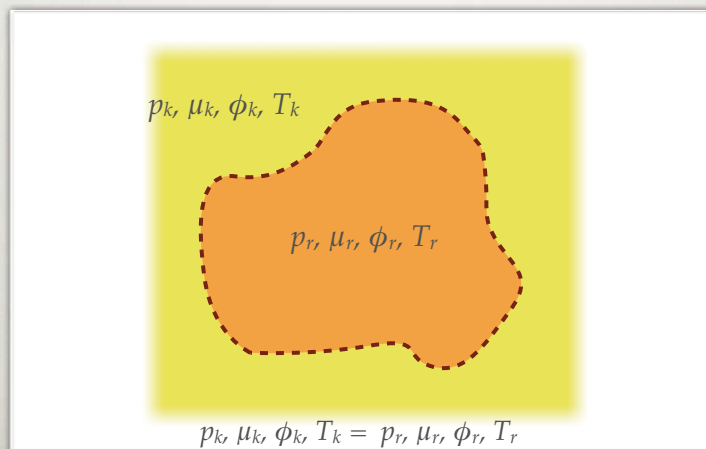
Change in the state functions depend only on the initial and final states of the system; it is independent of the path.



Other state functions: enthalpy ( $H$ ), Helmholtz free energy ( $F$ ), Gibbs free energy ( $G$ ), entropy ( $S$ )

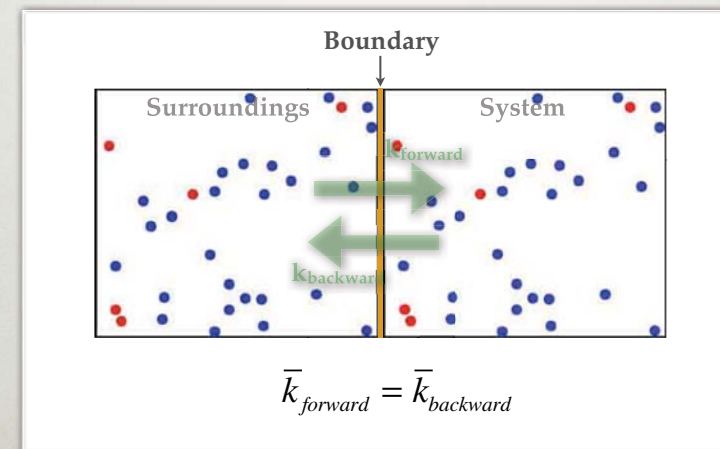
# EQUILIBRIUM

*Macroscopic description:* intensive variables are identical (between the system and surroundings, or between different parts of the system)



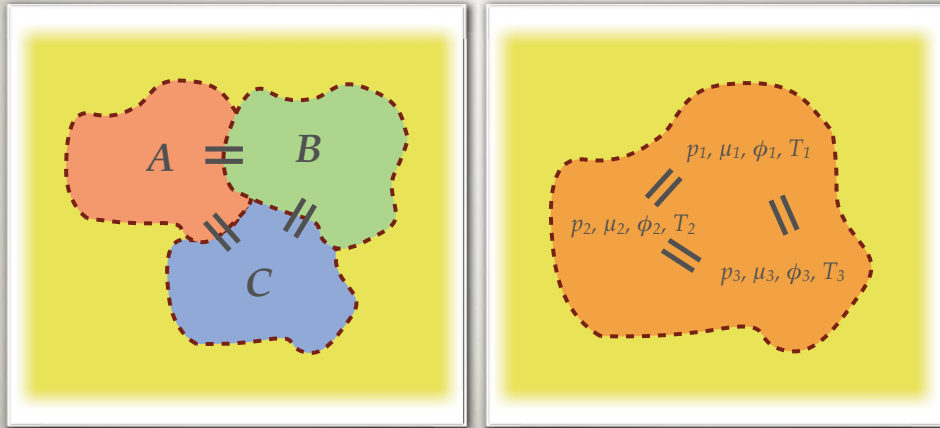
# EQUILIBRIUM

*Microscopic description:* average rates of forward and reverse reactions are equal.



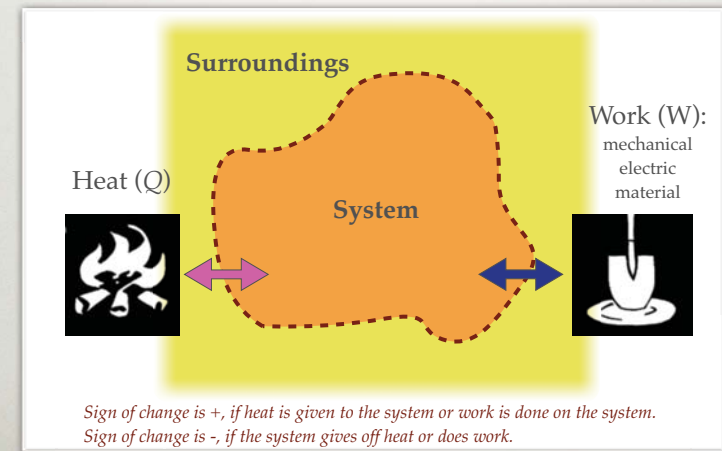
# ZEROth LAW OF THERMODYNAMICS

- If two systems (A and B) are independently in equilibrium with a third one (C), then they are in equilibrium with each other as well.
- Between different points of a system in equilibrium, the intensive variables are equal (there are no thermodynamic currents).



# CHANGE

Thermodynamics is interested in the changes within the system.  
Changes may be evoked by: heating, work.  
Result: internal energy of the system changes.

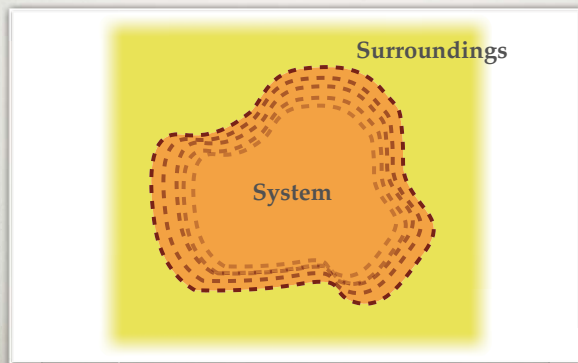


# EQUILIBRIUM CHANGE

Definitions:

- Reversible change (reversibly driven process).
- Series of reversible (equilibrium, infinitesimal) steps; the system is in equilibrium during each step of the process.
- "Infinitely slow", "quasi-static" process.
- There are no permanent (compensating) energy changes in the surroundings (work done by the system and surroundings are equal).

Characteristics: Work done by the system is maximal.



# THERMODYNAMIC PROCESSES

- **Isobaric:** at constant pressure.
- **Isochoric (isometric, isovolumetric):** at constant volume.
- **Isothermic:** at constant temperature.
- **Adiabatic:** no heat exchange.
- **Isentropic:** at constant entropy.
- **Isenthalpic:** at constant enthalpy.
- **"Steady state":** at constant internal energy.

# FIRST LAW OF THERMODYNAMICS

- Law of conservation of energy.
- Energy may be converted into different forms, but the total energy of the system remains constant.
- The change in the internal energy of the system is the sum of the supplied heat ( $Q_E$ ) and the work ( $W$ ) done on the system.

$$E_2 - E_1 = \Delta E$$

Difference between the internal energies of the initial and final states of the system.

$$\Delta E = Q_E + W$$

Work parts: *Volumetric work*  
 $W_V = -p\Delta V$   
 Negative, because in the case of work done on the system  $V_2 - V_1 < 0$ .

*Electric work:*  
 $W_Q = q\Delta Q$

*Material transport: Generalized:*  
 $W_n = \mu\Delta n$        $W^i = y_{int}^i \Delta x_{ext}^i$

Isochoric process ( $\Delta V=0$ ):

$$\Delta E = Q_V$$

Change in internal energy is caused only by the heat exchange.

Isobaric process ( $p=\text{constant}$ ):

$$\Delta E = Q_p + W_V = Q_p - p\Delta V$$

# ENTHALPY

- "Heat content" ( $H$ )
- Total heat absorbed by the system at constant pressure.
- Thermodynamic state function, by the help of which the useful work done at constant pressure may be calculated.
- *enthalpein* (Gr) to heat

$$\Delta E = Q_p + W_V = Q_p - p\Delta V$$

Internal energy change of the system during isobaric process.

Enthalpy:  $H = E + pV$

$E, p, V$ : equilibrium internal energy, pressure and volume.

Enthalpy change:  $\Delta H = \Delta E + p\Delta V$

Isobaric process:  $\Delta H = Q_p$

If there are no other energy changes beside volumetric work.  
 $Q_p$  = Reaction heat at constant pressure.

Isochoric process ( $V=\text{constant}$ ):  $\Delta E = Q_V$

Volumetric work = 0.  
 $Q_V$  = Reaction heat at constant volume.

Standard state: 25 °C, 1 atm. Standard enthalpy of compounds ( $\Delta H^\circ$ ): reaction heat of 1 mole of compound measured in standard state, at constant pressure.

# SECOND LAW OF THERMODYNAMICS

- "Law of entropy."
- Heat flows spontaneously from regions of high to low temperature.
- In an isolated system only such processes occur spontaneously which tend to equilibrate the respective intensive variables.
- During spontaneous processes entropy increases. Thermodynamic equilibrium is characterized by entropy maximum.
- Spontaneous processes proceed towards the most probable state (irreversible).
- Statistical law.



The fallen and broken flower pot does not spontaneously return to its original place.



The essentially endless heat content of oceans cannot be used to spontaneously produce useful work.

# ENTROPY

- *Entropy* ( $S$ ): unavailability of the system to do useful work. State function (units: J/K).
- Measure of "unusable" energy.
- Measure of thermodynamic probability.
- Measure of disorder.
- *entrepein* (Gr) to turn around.

## Phenomenologic description:

Energy change during thermodynamic processes (interactions):

$$\Delta E = \sum_i y_{int}^i \Delta x_{ext}^i$$

Product of intensive variables ( $y$ ) and extensive variable-changes ( $\Delta x$ ).

Energy change during thermal processes (thermal interactions):

$$Q_E = T\Delta S$$

Intensive variable: temperature ( $T$ )  
 Extensive variable: entropy change ( $\Delta S$ )

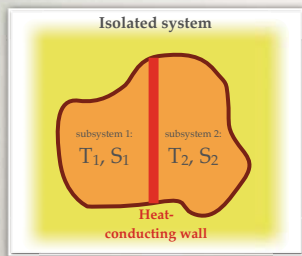
Entropy change:

$$\Delta S = \frac{\Delta Q}{T}$$

Measure of heat exchanged during reversible (equilibrium) isothermal process.

# ENTROPY

- In isolated system, during temperature equilibration entropy always increases.



$$\Delta E = \Delta(E_1 + E_2) = \Delta E_1 + \Delta E_2 = 0$$

In an isolated system internal energy is constant ( $\Delta E=0$ ).

$$\Delta E_1 = T_1 \Delta S_1 \quad \Delta E_2 = T_2 \Delta S_2$$

Energy changes during thermal interaction.

$$\Delta S = \Delta S_1 + \Delta S_2$$

Entropy is an extensive variable.

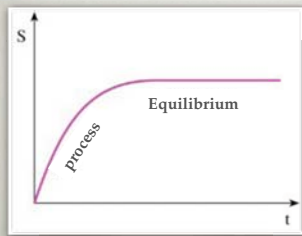
$$\Delta S = \frac{\Delta E_1}{T_1} + \frac{\Delta E_2}{T_2} = \Delta E_1 \left( \frac{1}{T_1} - \frac{1}{T_2} \right)$$

$$T_1 > T_2 \quad \text{or} \quad T_1 < T_2 : \Delta S > 0$$

If the initial temperatures are different, entropy increases.

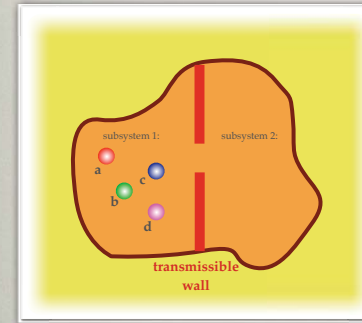
$$T_1 = T_2 : \Delta S = 0$$

If the initial temperatures are equal, entropy does not change.



# STATISTICAL INTERPRETATION OF ENTROPY

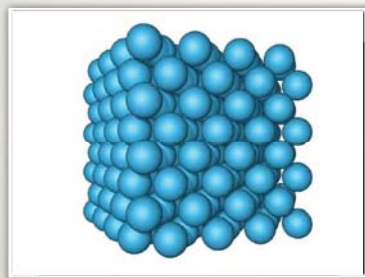
- *Macrostate*: state of the system defined by state variables (e.g.,  $p, V, T$ ).
- *Microstate*: internal structure of the system (position, velocity of every particle)
- In the *most probable macrostate* the number of microstates is the greatest.
- Number of microstates describing a given macrostate: *thermodynamic probability* ( $\Omega$ ).



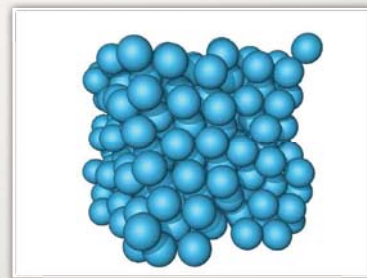
Macrostate	number of particles in subsystem 1	structure of subsystem 1	Number of microstates ( $\Omega$ )
A	4	abcd	1
B	3	abc,abd,acd,bcd	4
C	2	ab, ac, ad, bc, bd, cd	6
D	1	a, b, c, d	4
E	0	-	1

Entropy:  $S = k_B \ln \Omega$      $k_B$  = Boltzmann's constant ( $1,38 \times 10^{-23} \text{ J/K}$ )

# ENTROPY IS THE MEASURE OF DISORDER



Ordered crystal:  
Low entropy



Disordered liquid:  
High entropy

# THIRD LAW OF THERMODYNAMICS

- Nernst's law.
- The entropy of mon-component, crystallizing material at 0 K temperature is 0.

At 0 K thermal motions halt, there are no thermal crystal faults.

$$\Omega = 1$$

Number of microstates is 1.

Because of mono-component system, only one type of molecular arrangement is possible.

$$S = k_B \ln 1 = 0$$

# GIBBS FREE ENERGY

- Gibbs free energy (G), Gibbs potential.
- Useful work extractable from isothermic, isobaric (not isolated) system.
- State function

Entropy is an extensive variable.

$S_{tot}$  = total entropy (entropy of the universe)  
 $S_{sys}$  = system entropy  
 $S_{sur}$  = entropy of surroundings

$$\Delta S_{tot} = \Delta S_{sys} + \Delta S_{sur}$$

$$\Delta S_{sur} = -\frac{\Delta H_{sys}}{T}$$

$$\Delta S_{tot} = \Delta S_{sys} - \frac{\Delta H_{sys}}{T}$$

$$-T\Delta S_{tot} = \Delta H_{sys} - T\Delta S_{sys}$$

$$\Delta G = \Delta H_{sys} - T\Delta S_{sys}$$

$$H = G + TS$$

Enthalpy = useful work + unusable energy change

During spontaneous processes the entropy of the universe increases ( $\Delta S_{tot} > 0$ ):

$$\Delta S_{sys} > \frac{\Delta H_{sys}}{T}$$

During spontaneous processes, Gibbs free energy decreases ( $\Delta G < 0$ ):

$$\Delta G = \Delta H_{sys} - T\Delta S_{sys} < 0$$

# HELMHOLTZ FREE ENERGY

- Helmholtz free energy (F)
- Useful work extractable from isothermic, isochoric (constant volume) (not isolated) system.
- State function

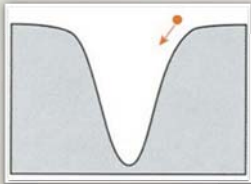
$$F = E - TS$$

F : Helmholtz free energy  
 E: internal energy  
 T: absolute temperature  
 S: entropy

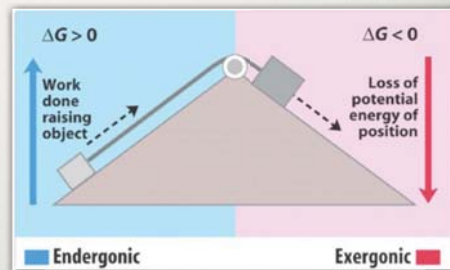
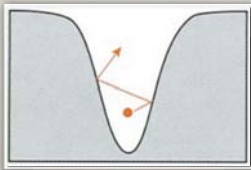
## SIGNIFICANCE OF ENTROPY AND GIBBS FREE ENERGY

- Spontaneous processes are those, during which Gibbs free energy decreases ( $\Delta G < 0$ ) and the entropy of the universe increases ( $S_{tot} > 0$ ).

Decrease in free energy: stabilizes



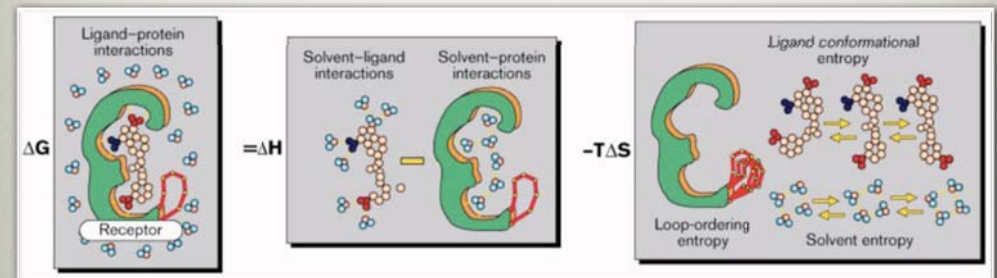
Increase in entropy: randomizes



- Useful work is done by exergonic processes.
- Endergonic processes can be driven by coupling to exergonic processes.
- Entropy can be decreased locally.
- Life consumes entropy (its entropy is decreased at the expense of increasing total entropy).

## SIGNIFICANCE OF ENTROPY AND GIBBS FREE ENERGY

- Enthalpic component: molecular interactions (covalent, van der Waals, electrostatic, hydrogen bond)
- Entropic component: internal structural disorder (arrangement of solvent molecules and counter-ions, degrees of freedom, rotational and translational arrangements)



Gibbs free-energy change associated with receptor-ligand interaction binding

Enthalpic component: changes in molecular interactions

Entropic component: conformational changes

# CHEMICAL POTENTIAL

- Relevant biomedical thermodynamic systems are usually mixtures of chemicals.
- Chemical potential characterizes the thermodynamics of mixtures.
- Mixtures contain many components in different concentrations.

Gibbs free energy of a mixture: sum of free energies of components

$$G = \sum_{i=1}^k n_i \mu_i$$

$k$  = number of components  
 $n_i$  = amount of  $i^{\text{th}}$  component (mol)  
 $\mu_i$  = free energy of one mol of the  $i^{\text{th}}$  component (**chemical potential**)

$$\mu_i^0$$

**standard chemical potential** of the  $i^{\text{th}}$  component:  
 molar free energy at standard conditions  
 (25 °C, 1 atm, 1 mol/l)

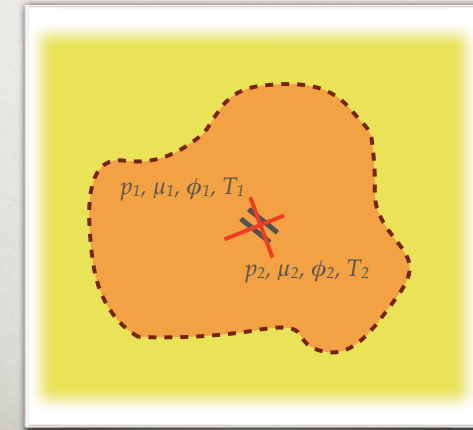
*Meaning: free-energy change associated with the addition/removal of 1 mol of the  $i^{\text{th}}$  component into/from the system.*

$$\mu_i = \mu_i^0 + RT \ln \frac{c_i}{c_i^0}$$

$c_i$  = concentration of the  $i^{\text{th}}$  component (mol/l)  
 $c_i^0$  = 1 mol/l

# IRREVERSIBLE THERMODYNAMICS

- Natural processes are rarely reversible.
- If there are inequalities in the intensive variables at different locations within the system, thermodynamic currents arise.
- Thermodynamic currents aim at the restoration of equilibrium.



# THERMODYNAMIC CURRENTS

- Natural processes are rarely reversible.
- If there are inequalities in the intensive variables at different locations within the system, thermodynamic currents arise.
- Thermodynamic currents aim at the restoration of equilibrium.

Thermodynamic current	Relevant intensive variable (its difference maintains current)	Current density	Physical law
Heat flow	Temperature ( $T$ )	$J_E = -\lambda \frac{\Delta T}{\Delta x}$	Fourier
Volumetric flow	Pressure ( $p$ )	$J_v = -\frac{R^2}{8\eta} \frac{\Delta p}{\Delta x}$	Hagen-Poiseuille
Electric current	Electric potential ( $\phi$ )	$J_Q = -\frac{1}{\rho} \frac{\Delta \phi}{\Delta x}$	Ohm
Material transport (diffusion)	Chemical potential ( $\mu$ )	$J_n = -D \frac{\Delta c}{\Delta x}$	Fick

# HEAT FLOW

Thermodynamic current	Relevant intensive variable (its difference maintains current)	Current density	Physical law
Heat flow	Temperature ( $T$ )	$J_E = -\lambda \frac{\Delta T}{\Delta x}$	Fourier

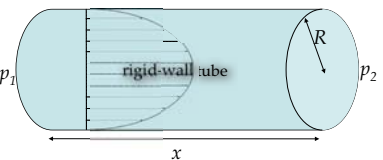
$Q$  = heat  
 $t$  = time  
 $R$  = tube radius  
 $x$  = length  
 $(\Delta T/\Delta x)$  = temperature gradient, maintained by  $T_1 - T_2$   
 $A$  = cross-sectional area of tube  
 $J_E$  = heat current density  
 $\lambda$  = coefficient of heat conductance

$$\frac{Q}{At} = J_E = -\lambda \frac{\Delta T}{\Delta x}$$

$$A = R^2 \pi$$

# VOLUME FLOW

Thermodynamic current	Relevant intensive variable (its difference maintains current)	Current density	Physical law
Volumetric flow	Pressure ( $p$ )	$J_V = -\frac{R^2 \Delta p}{8\eta \Delta x}$	Hagen-Poiseuille



$V$  = volume  
 $t$  = time  
 $R$  = tube radius  
 $\eta$  = viscosity  
 $p$  = pressure  
 $x$  = tube length  
 $(\Delta p/\Delta x)$  = pressure gradient, maintained by  $p_1$ - $p_2$ )  
 $A$  = cross-sectional area of tube  
 $J_V$  = flow intensity

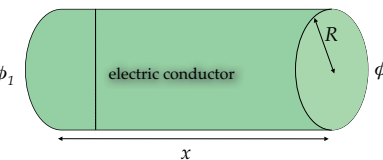
$$\frac{V}{t} = \frac{R^4 \pi \Delta p}{8\eta \Delta x}$$

$$A = R^2 \pi$$

$$J_V = \frac{V}{tA} = \frac{R^2 \Delta p}{8\eta \Delta x}$$

# ELECTRIC CURRENT

Thermodynamic current	Relevant intensive variable (its difference maintains current)	Current density	Physical law
Electric current	Electric potential ( $\phi$ )	$J_Q = -\frac{1}{\rho} \frac{\Delta \phi}{\Delta x}$	Ohm

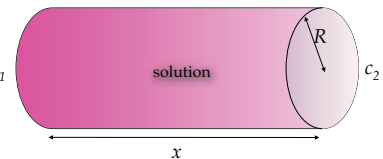


$q$  = electric charge  
 $t$  = time  
 $R$  = tube radius  
 $\phi$  = electric potential  
 $x$  = length of conductor  
 $(\Delta \phi/\Delta x)$  = potential gradient (voltage), maintained by  $\phi_1$ - $\phi_2$ )  
 $A$  = cross-sectional area of tube  
 $J_Q$  = electric current

$$\frac{q}{tA} = J_Q = -\frac{1}{\rho} \frac{\Delta \phi}{\Delta x}$$

# MATERIAL TRANSPORT (DIFFUSION)

Thermodynamic current	Relevant intensive variable (its difference maintains current)	Current density	Physical law
Material transport (diffusion)	Chemical potential ( $\mu$ )	$J_n = -D \frac{\Delta c}{\Delta x}$	Fick



$m$  = amount of material transported  
 $t$  = time  
 $R$  = tube radius  
 $x$  = length  
 $(\Delta c/\Delta x)$  = concentration gradient, maintained by  $c_1$ - $c_2$ )  
 $A$  = cross-sectional area of tube  
 $J_n$  = heat current density  
 $D$  = diffusion coefficient

$$\frac{m}{tA} = J_n = -D \frac{\Delta c}{\Delta x}$$

# THERMODYNAMIC CURRENTS

- Natural processes are rarely reversible.
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Electric current	Electric potential ( $\phi$ )	$J_Q = -\frac{1}{\rho} \frac{\Delta \phi}{\Delta x}$	Ohm
Material transport (diffusion)	Chemical potential ( $\mu$ )	$J_n = -D \frac{\Delta c}{\Delta x}$	Fick

# ONSAGER'S LINEAR FUNCTION

- Lars Onsager (1903-1976), Nobel prize (1968)
- General relationship between transport processes.

$$J = LX$$

Flow intensity of transported extensive variable (thermodynamic current,  $J$ )

$$J = \frac{\Delta x_{ext}}{A \Delta t}$$

=

Coefficient of conductance (thermodynamic coefficient,  $L$ )

x

Change in intensive variable ("gradient", thermodynamic force,  $X$ )

$$X = -\frac{\Delta y_{int}}{\Delta x}$$

A thermodynamic current may be evoked by several different intensive variables.  
E.g., thermodiffusion (material transport evoked by temperature gradient).

$$J_1 = L_{11}X_1 + L_{12}X_2$$

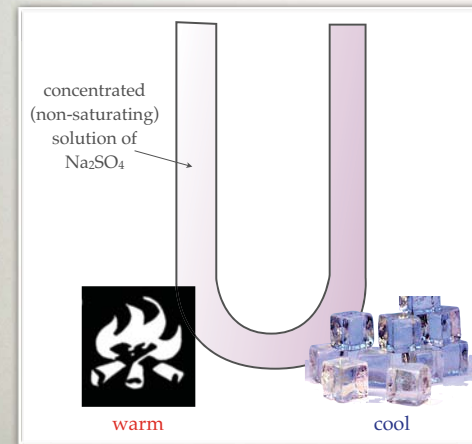
$$J_2 = L_{21}X_1 + L_{22}X_2$$

$L_{11}$ ,  $L_{22}$  = "linear" coefficients of conductance

$L_{12}$ ,  $L_{21}$  = "cross" coefficients of conductance

# THERMODIFFUSION

- Ludwig-Soret effect
- Carl Friedrich Wilhelm Ludwig (1856)
- Charles Soret (1879)



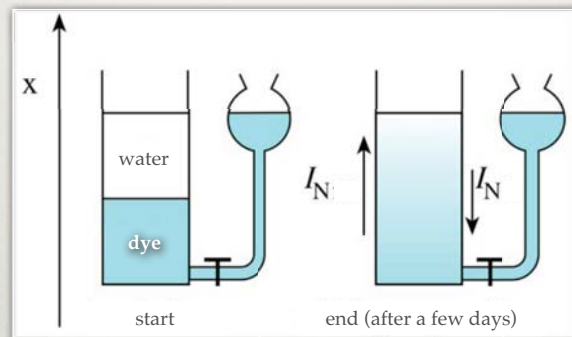
Two types of currents:

material flow  $J_V = -L_T \frac{\Delta T}{\Delta x}$

heat flow flow  $J_E = -\lambda \frac{\Delta T}{\Delta x}$

# DIFFUSION

- Spontaneous mixing, distribution, concentration-equilibration driven by the thermal motion of particles.



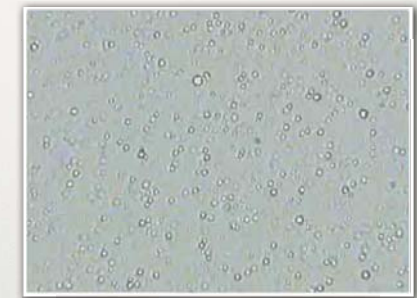
$$x^2 = 2Dt$$

$x$  = displacement of boundary  
 $t$  = time  
 $D$  = constant ("diffusion coefficient")

# Microscopic manifestation of diffusion: Brownian movement

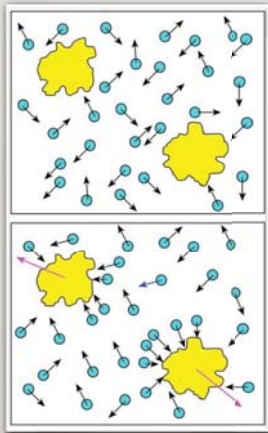


Robert Brown  
(1773-1858)

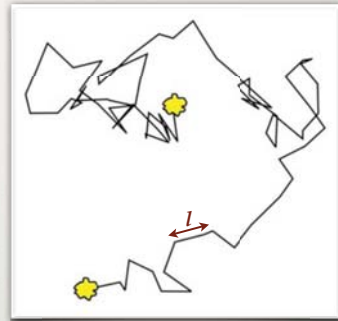


Fat droplets suspended in milk (droplet sizes 0.5 - 3  $\mu\text{m}$ )

# Brownian movement



Movement of microscopic particle is the result of random collisions with molecules.



$l$  = mean free path (average distance between consecutive collisions)  
 $v$  = average velocity of thermally propelled particle

# DIFFUSION

- *Fick's First law*: material flow density is the product of the eviking concentration gradient and the diffusion coefficient

Material transport:  $J_n = -D \frac{\Delta c}{\Delta x}$

$J_n$  = material flow  
 $\Delta c / \Delta x$  = drop in concentration ("gradient")  
 $D$  = constant ("diffusion coefficient")

Diffusion coefficient:  $D = \frac{1}{3} v l$

$v$  = average velocity of thermally propelled particle  
 $l$  = mean free path (average distance between consecutive collisions)  
 $D$  = amount of material transported through unit area per unit time ( $\text{m}^2/\text{s}$ ) (at unitary concentration drop).



Brown-movement

Diffusion coefficient for spherical particle:  $D = \frac{k_B T}{6 \pi \eta r}$

*Einstein-Stokes formula:*  
 $k_B$  = Boltzmann's constant  
 $T$  = absolute temperature  
 $\eta$  = viscosity  
 $r$  = radius of particle

# DIFFUSION

- *Fick's Second Law*: instantaneous material flow density depends on the temporal change in evoking concentration gradient.

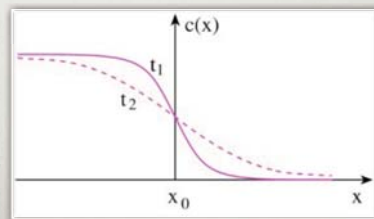
Material flow intensity:  $-\frac{\Delta J_n}{\Delta x} = \frac{\Delta c}{\Delta t}$

$J_n$  = Flow intensity  
 $x$  = distance  
 $t$  = time

Diffusion coefficient:  $D \frac{\Delta \left( \frac{\Delta c}{\Delta x} \right)}{\Delta x} = \frac{\Delta c}{\Delta t}$

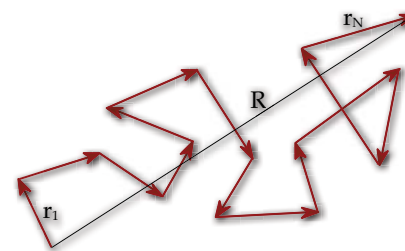
$D$  = diffusion coefficient

Concentration gradient decreases with time (boundary becomes "smeared")



# RELATIONSHIP BETWEEN RANDOM WALK AND DIFFUSION

## Brown-movement - "random walk"



Mean particle velocity:  $v = \frac{l}{\tau}$

Full walk time:  $t = N \tau$

Diffusion coefficient:  $D = \frac{1}{3} v l$

## "Square-root law":

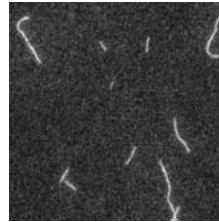
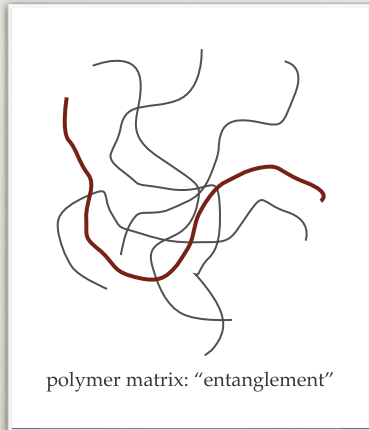
$$\langle R^2 \rangle = N l^2 = L l$$

$R$  = displacement  
 $N$  = number of unit vectors (steps)  
 $l = |\vec{r}_i|$  = mean free path  
 $r_i$  = unit step  
 $Nl = L$  = total distance

$$\langle R \rangle = \sqrt{N l^2} = \sqrt{\frac{t}{\tau} l^2} = \sqrt{t v l} = \sqrt{3 D t}$$

## A SPECIAL CASE OF DIFFUSION: REPTATION

- *Reptation*: “snake-like” diffusion within polymer network (*Reptilia*)



Actin filaments in methyl-cellulose matrix.  
“unidirectional diffusion”

## A SPECIAL CASE OF DIFFUSION: OSMOSIS

**Osmosis** ~ osmos (Gr) thrust, push

**History:** -Nollet (1748): pig bladder containing alcohol swells in water  
-Dutrochet (1830): first osmometer  
-van't Hoff: gas-law mechanism of osmosis  
-Arrhenius and Findlay: vapor-pressure mechanism of osmosis

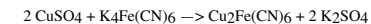
**Osmosis:** diffusion-driven process between two compartments separated by semipermeable membrane (solvent diffusion).

**Osmotic pressure ( $\pi$ )** = equilibrium hydrostatic pressure between each side of the semipermeable membrane (dP).

Semipermeable membranes:

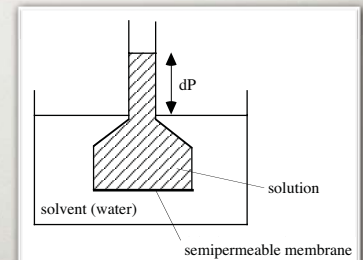
-frog membrane

-copper-ferrocyanide gel (precipitated on the pores of clay cylinder):



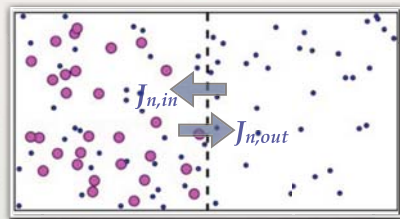
-cellulose membrane, dialysis bag

Biomedical importance: biological membranes, compartment surfaces.



## A SPECIAL CASE OF DIFFUSION: OSMOSIS

- *Osmosis*: one-directional solvent transport by diffusion.



$J_{n,in}$  driving force: solvent concentration difference

$J_{n,out}$  driving force: pressure difference

Osmotic equilibrium:  $J_{n,in} = J_{n,out}$

Osmotic pressure:  $p_{osmosis} = cRT$

Importance of osmosis: cell swelling, edema, hemodialysis.

*van't Hoff's Law:*  
 $c$  = concentration of solute  
 $R$  = universal gas constant  
 $T$  = absolute temperature

## MECHANISMS OF OSMOSIS I.

*van't Hoff's gas-law mechanism*

Universal gas law:  $pV = RT$

$$p = \frac{1}{V}RT$$

Osmotic pressure:  $p_{osm} = \pi = cRT$

$p$  = pressure  
 $V$  = volume  
 $R$  = gas constant (8,3 J/mol.K)  
 $T$  = absolute temperature  
 $c$  = molar concentration of solute

**Problem:**

What is the osmotic pressure of a 0.1 M (0.1 mol/dm<sup>3</sup>) sucrose solution?

$$\pi = 8.3 \text{ (J/mol.K)} \times 293 \text{ (K)} \times 0.1 \text{ (mol/dm}^3\text{)} = 243 \text{ kPa} \sim 2.4 \text{ atm.}$$

Note: the formula applies to dilute solutions.

# MECHANISMS OF OSMOSIS II.

## Corrected van't Hoff's law

$$\pi = fnRTc$$

$n$  = number of ions produced during complete dissociation

$f$  = osmotic coefficient:

$$f = \frac{\pi_{measured}}{\pi_{calculated}} = \frac{\pi_{measured}}{nRTc}$$

## Arrhenius-Findlay's vapor pressure mechanism

Changes in the chemical potential of the solvent due to the solute (chemical activity of the solvent becomes reduced upon the addition of the solute).

$$\text{Chemical potential of solvent: } \mu_s = \mu_s^0 + RT \ln \chi_s$$

$\mu_s^0$  = standard chemical potential of solvent  
 $\chi_s$  = mole fraction of solvent

$$\text{Mole fraction of solvent: } \frac{n_0}{n_0 + n_1}$$

$n_0$  = number of moles of solvent  
 $n_1$  = number of moles of solute

$$\text{Raoult's law: } \frac{p_1}{p_0} = \frac{n_0}{n_0 + n_1}$$

$p_0$  = vapor pressure of pure water  
 $p_1$  = vapor pressure of solution water

$$\text{Arrhenius-Findlay law: } \pi = \frac{RT}{v} \ln \frac{p_0}{p_1}$$

$v$  = molar volume of solvent (cm<sup>3</sup>/mol; for water: 18)

# IMPORTANCE OF OSMOSIS

Properties of the semipermeable membrane: **reflexion coefficient** ( $\sigma$ )

For a perfect semipermeable membrane:

$$\Delta P = \Delta \pi \quad \text{and} \quad \sigma = \frac{\Delta P}{\Delta \pi} = 1$$

In reality:

$$\Delta P < \Delta \pi \quad \text{and} \quad \sigma = \frac{\Delta P}{\Delta \pi} < 1$$

$$0 < \sigma < 1$$

**Osmotic work:**

$$-L = nRT \ln \frac{c_1}{c_2} = nRT \ln \frac{\pi_1}{\pi_2}$$

$R$  = gas constant

$n$  = total number of moles of solute

$c_1$  = initial molarity of solution

$c_2$  = final molarity of solution (after equilibration)

$\pi_1$  = initial osmotic pressure

$\pi_2$  = final osmotic pressure (after equilibration)

**Equivalent osmotic pressure** (osmotic concentration): concentration of non-electrolyte solution in equilibrium with a heterogenous solution system.

Units: mmol/kg = mOsmol/kg = mOsm

Osmolarity of blood plasma ~ 300 mOsm.

**Oncotic pressure:** colloid osmotic pressure. Osmotic pressure of a solution of colloidal macromolecules.

