

Condensed matter structure and properties

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Febr.9, 2010

1. Solid state materials - structure

crystalline

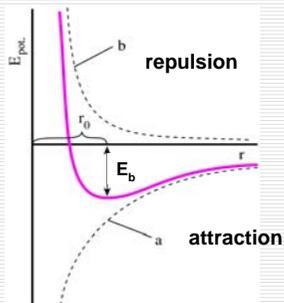
or

(amorphous)

1. Solid state materials - structure

Structure of crystalline materials

Ideal crystals: characterized by order



$$E_{pot} = \frac{B}{r^m}$$

$$m > n$$

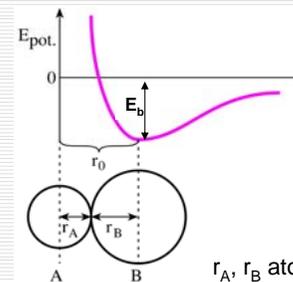
$$E_{pot} = -\frac{A}{r^n}$$



constituting atoms are at bond distances from each other

r = distance of two atoms forming the bond
 r_0 = bond distance
 E_{bond} = bond energy

Structure of crystalline materials



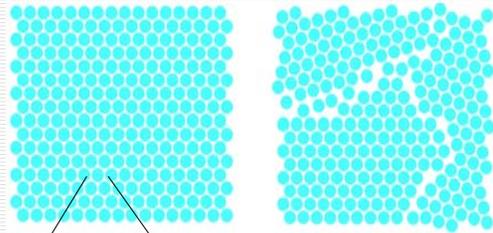
r_A, r_B atomic radii of atoms A and B

Atom	Z	van der Waals	Co-valent	Ionic	Ion
H	1	0,120	0,037	-	H ⁺
C	6	0,170	0,077	0,029	C ⁺
N	7	0,155	0,075	0,025	N ⁺
O	8	0,152	0,073	0,140	O ²⁻
F	9	0,147	0,071	0,117	F ⁻
P	15	0,180	0,106	0,058	P ³⁺
S	16	0,180	0,102	0,184	S ²⁻

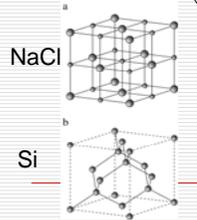
Atomic/molecular bond distances ($r_0 = r_A + r_B$) and binding energies E_b depend on the type of interactions (functions **a** and **b**)

Structure of crystalline materials

ideal crystal – single crystal

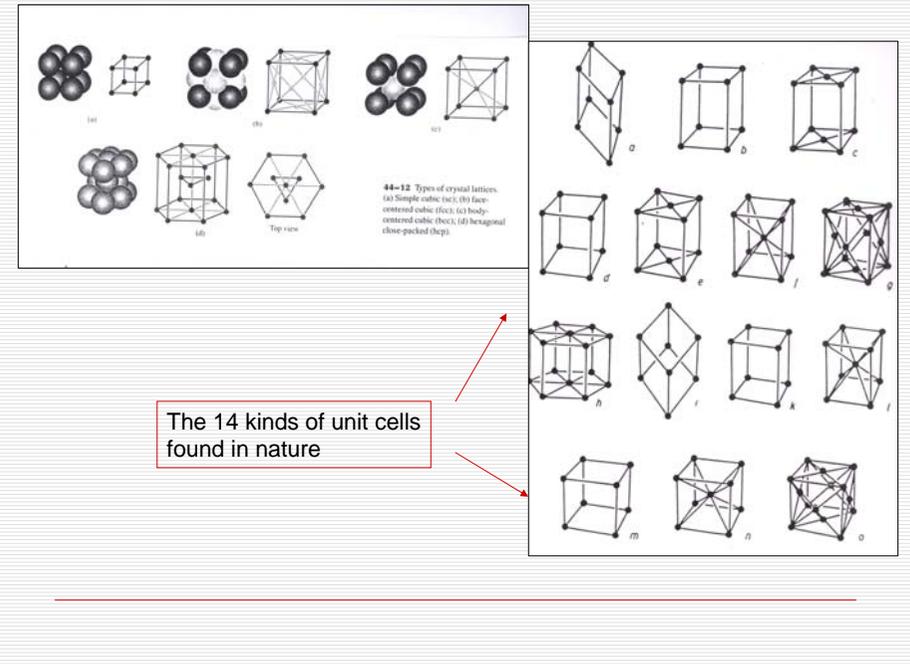


microcrystalline state
(real crystals)



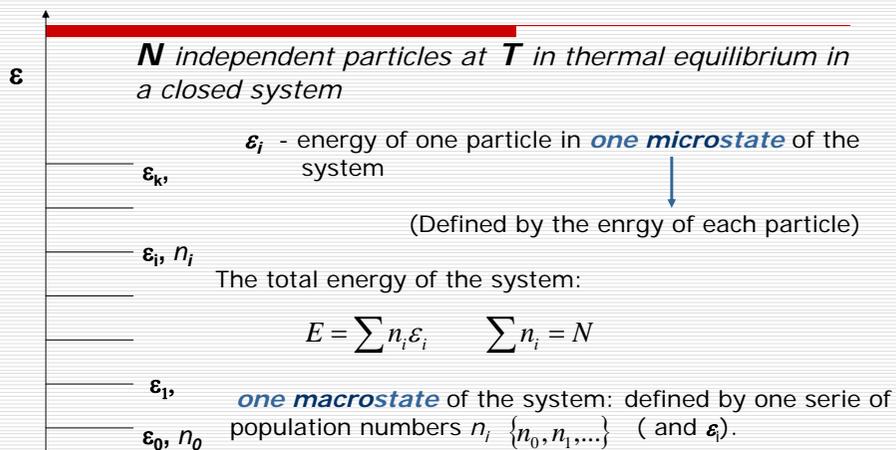
unit cell
is repeated
to build up the
crystal lattice

14 kinds of unit cells
are found in nature



The 14 kinds of unit cells
found in nature

Reminder: Boltzmann distribution



Reminder: Boltzmann distribution

Boltzmann distribution – important formulae

$$p_i = \frac{e^{-\frac{\epsilon_i}{kT}}}{Z}$$

probability that particles exist with ϵ_i energy in the system

$$Z = \sum_i e^{-\frac{\epsilon_i}{kT}}$$

partition function

$$n_i = \frac{N e^{-\frac{\epsilon_i}{kT}}}{Z}$$

$$\frac{n_i}{n_0} = e^{-\frac{\epsilon_i - \epsilon_0}{kT}} = e^{-\frac{\Delta\epsilon}{kT}}$$

Boltzmann factor

The relative population of two energy states

Structure of crystalline materials

Considering two states of bonds – intact
- broken

$$\frac{n_b}{n_i} = e^{-\frac{E_b}{kT}} \approx \frac{n}{N}$$

Boltzmann distribution

$$n_i + n_b = N$$

e.g. NaCl, $E_b=6.31\text{eV}$, $kT(\text{room tempr.})=0.023\text{eV}$
Atomic cohesive energy

$$\frac{n_b}{n_i} = e^{-\frac{6.31}{0.023}} \approx e^{-274} \approx 10^{-110} = 0$$

$E_b \sim$ primary bonds : covalent ↑ E_b
ionic
metallic

1 - 13 eV $E_b \gg kT$ (RoomTemp)

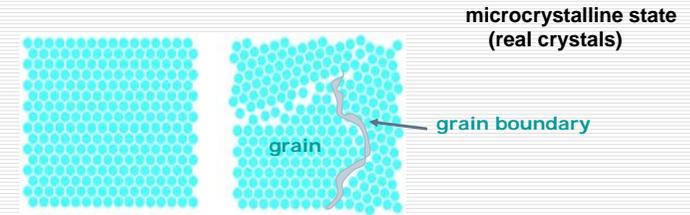
There would be perfect order in ideal crystals – no broken bonds

Structure of crystalline materials

primary bonds -> strong bonds -> not broken thermally at rt

long range order in ideal crystals: large number of atoms in periodic array

Still: crystal defects in real crystals

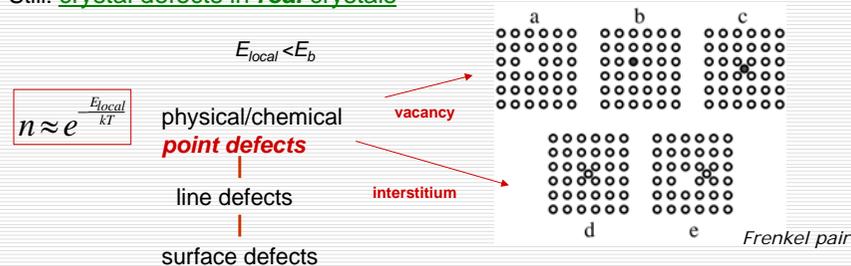


Structure of crystalline materials

primary bonds -> strong bonds -> not broken thermally at rt

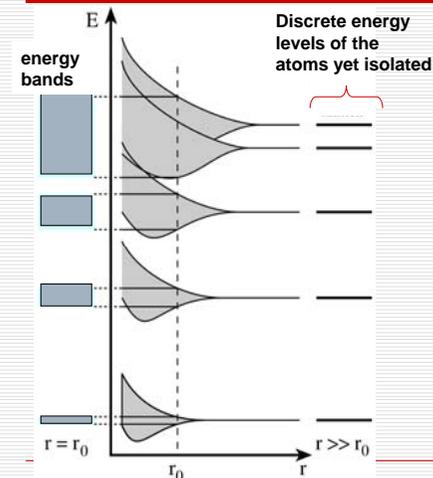
long range order in ideal crystals: large number of atoms in periodic array

Still: crystal defects in real crystals



2. Solid state materials (crystals)- properties

Interaction of ordered atoms changes the electronic energy levels



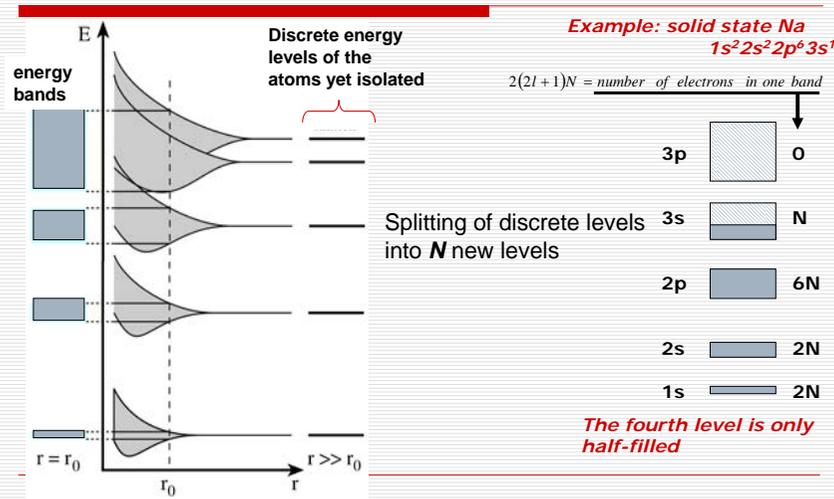
$N \sim 10^{23}$

N no. of isolated atoms with discrete energy levels change their electronic states when they interact to form a crystalline state with bond distance r_0 .

Splitting of discrete levels into N new levels results in continuous ranges :
energy bands

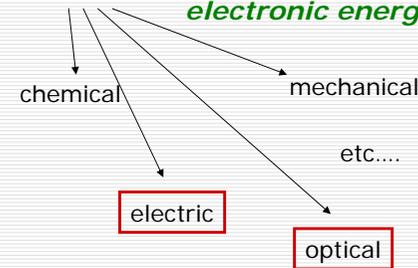
2. Solid state materials (crystals)- properties

Interaction of ordered atoms changes the electronic energy levels

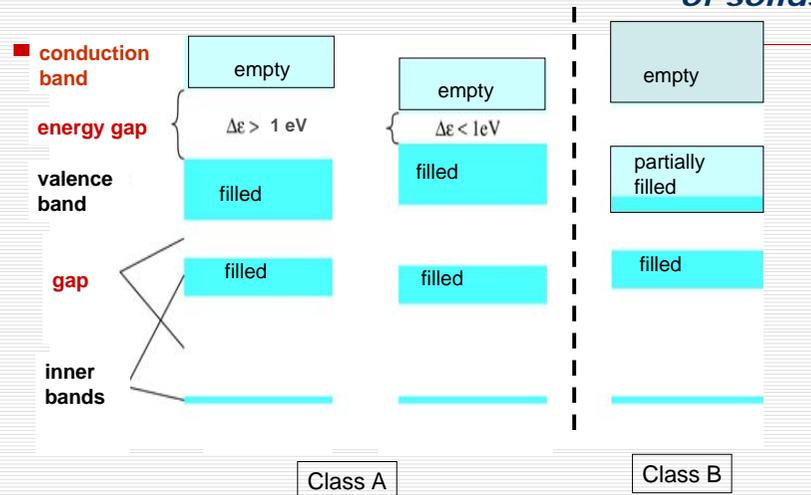


Significance of electronic energy level structure

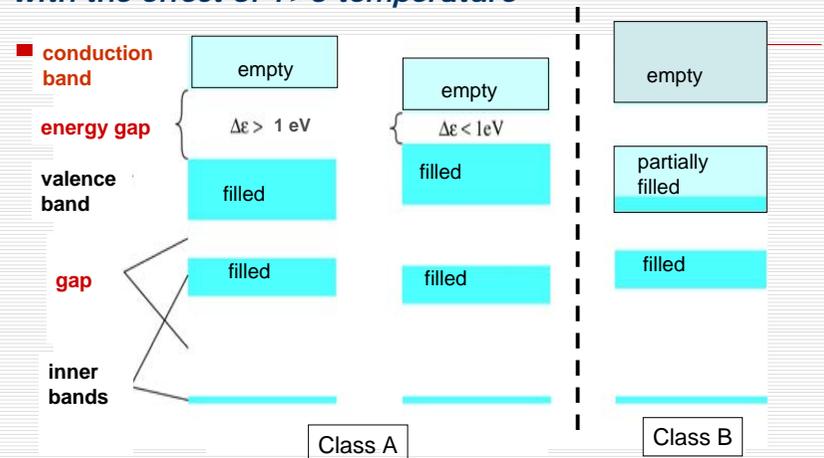
Properties of materials are the consequence of the **electronic energy levels**



4. Possible classes of electronic band features of solids

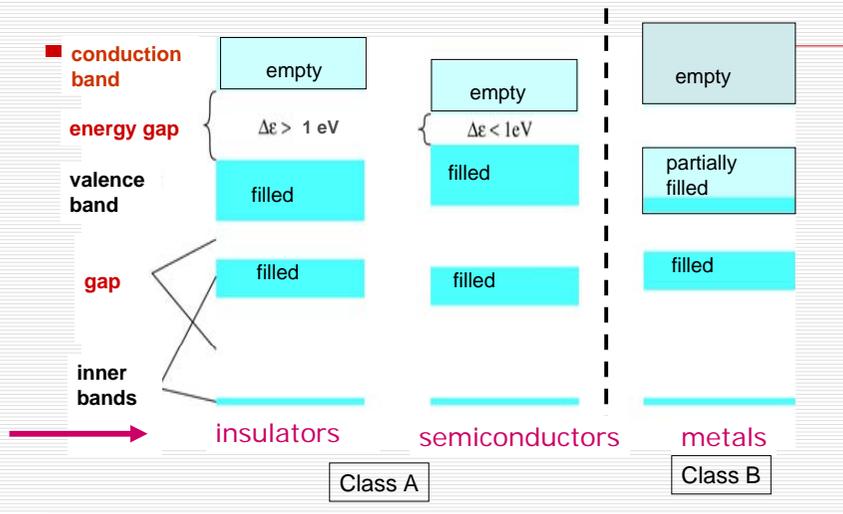


The schematics of band structures must be completed with the effect of $T > 0$ temperature

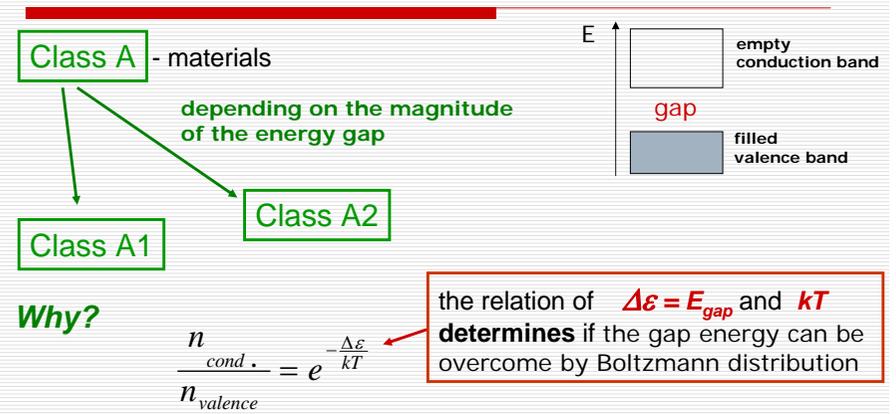


The concept of Boltzmann distribution is also applicable for the population of the electronic energy levels

5. Properties related to the energy band feature in solids



Classes of crystalline solid materials – band features and properties



Class A1

$$E_{gap} \gg 1 \text{ eV}$$

insulators

e.g. diamond $E_{gap} = 5.4 \text{ eV}$

$$\frac{N_{cond}}{N_{val}} = e^{-\frac{5.4}{0.023}} = e^{-235} \approx 0$$

- No electric conductivity (dielectric break-down: $\sim V/\text{bond} \rightarrow 10^{10} \text{ V/m}$)
- No photon absorption in the **VIS range** \rightarrow **transparency**
- UV photons may be absorbed, IR: excitation of lattice vibrations

Class A2

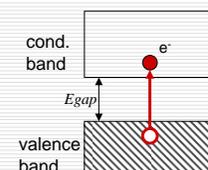
$$E_{gap} \leq 1 \text{ eV}$$

intrinsic semiconductors

Non-transparent crystalline materials

Reasonable number of thermally excited electrons in the conduction band

	$E_g \text{ (eV)}$	$\frac{N_{cond}}{N_{val}} = e^{-\frac{0.75}{0.023}} = e^{-33} = 7 \cdot 10^{-15}$
Si	1.1	
Ge	0.75	$N_{val} \approx 6 \cdot 10^{23} \Rightarrow N_{cond} \approx 4 \cdot 10^8 \text{ /cm}^3$



n - type conductivity (electron conduction)
 $\sigma \approx e \frac{E_{gap}}{2kT}$

p - type conductivity (electron-hole: + charge conduction)

Two kinds of charge carriers

Class A2

intrinsic semiconductors

$$\sigma = \text{const} * e^{-\frac{E_{gap}}{2kT}}$$

Slightly depends on T

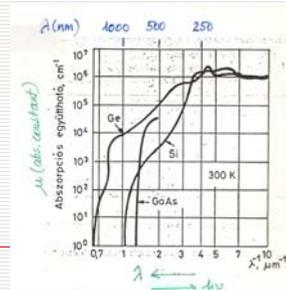
Equilibrium: generation and recombination of charge carriers is of equal probability
 $p(\text{recombination}) \sim n^2$, $p(\text{generation}) \sim \text{Boltzmann factor}$

Specific conductivity is increased by temperature increase
 → **thermistors**

Optical properties: non-transparency in the VIS range

$$hf_{VIS} > E_{gap}$$

Photon absorption induces conductivity
 → **photodetectors**



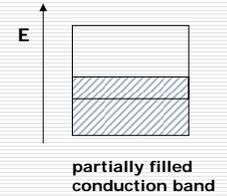
Class B

metals

e.g. 1-valence and 2-valence metals Na, Mg, Cu..

	Cu	Si	T=293 K
n(charge)/m ³	9x10 ²⁸	1x10 ¹⁶	
specific resistance (Ohmxm)	2x10 ⁻⁸	3x10 ³	

↑
high electric conductivity



Energy absorption is possible within the partially filled highest energy band

- Electrons conduct electricity
- Optical non-transparency

$$\sigma \approx \frac{1}{T}$$

Specific conductivity decreases with T-increase

↙ **semiconductors**

Class B

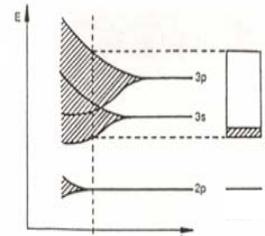
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n(charge)/m ³	9x10 ²⁸	1x10 ¹⁶
specific resistance (Ohmxm)	2x10 ⁻⁸	3x10 ³

↑
high electric conductivity

The example of Na-metal
 $1s^2 2s^2 2p^6 3s^1$



Broadening → overlap of bands 3s and 3p
 1/8-th filling of last band

Energy absorption is possible within the partially filled highest energy band

Class A2**

doped semiconductors

Doping: incorporation of a second component (dopant) into the crystal lattice of an intrinsic semiconductor (host) in a small amount

$$\frac{N_{host}}{N_{dopant}} \approx 10^6$$

→ **Dopant atoms are isolated in the crystal matrix**

Idea: properly selected **dopant may reduce E_{gap}** , thus increasing the number of thermally excited charge carriers : electrons or electron-holes

Class A2**

doped semiconductors

$$\frac{N_{\text{host}}}{N_{\text{dopant}}} \approx 10^6$$

Dopants are isolated in the crystal matrix

Dopants reduce the effective E_{gap}

Two combinations:

-4-valent host combined with 5-valent dopant → *n-type* doped semiconductor

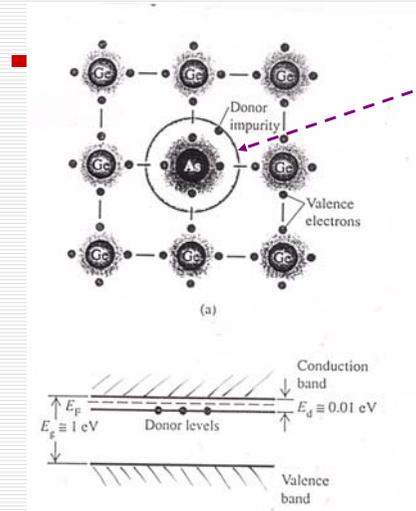
-4-valent host combined with 3-valent dopant → *p-type* doped semiconductor

Hosts: Ge, Si

Dopants: 5-valent : P, As, Bi

3-valent : B, Al, Ga, In

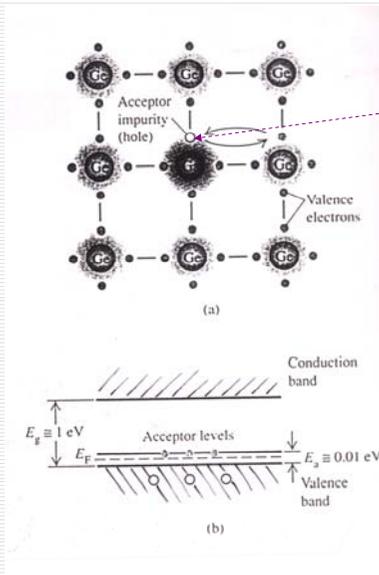
Example: Doping 4 valence Ge crystal with 5 valence As



Loosely bound electron, easy to excite into the delocalized conduction band → *n-type* conductivity

Donor level exists only at the isolated impurity atoms

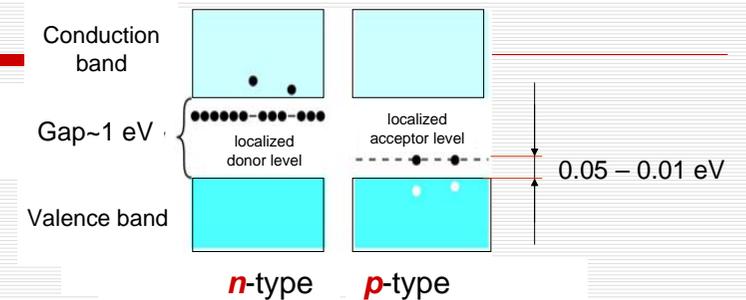
Example: Doping 4 valence Ge crystal with 3 valence Ga



One electron is missing at Ga to form covalent bond with one Ge neighbor → It can easily bind electrons from Ge-s → *p-type* conductivity

Acceptor level exists only at the isolated impurity atoms, but the electron-hole can diffuse among the Ge-valence electrons

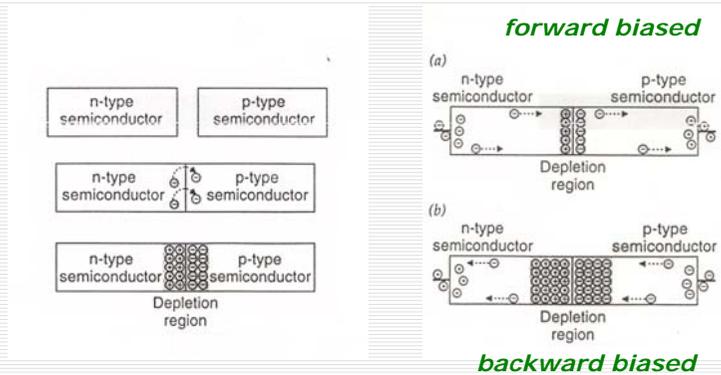
Summary: *n-type* and *p-type* conduction



electrons thermally excited from the donor level of impurity conduct electricity

electrons thermally excited to the acceptor level of impurity are localized, but *electron-holes* conduct electricity

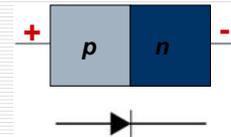
p – n junctions → revolution in electronics



diode and transistor:

constructed from n- and p- type doped semiconductors

forward biased

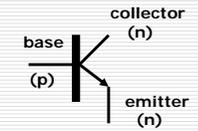
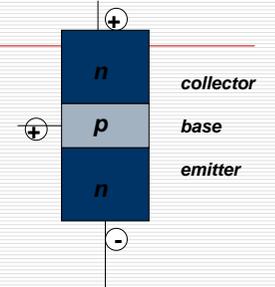


diode:

- rectifier
- light source under voltage
- transducer of light into voltage

transistor:

- electric current amplifier
- digital unit of computers



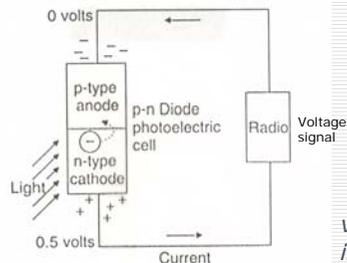
Doping makes possible miniature sizes → microelectronics

p – n junctions : light emitters and detectors

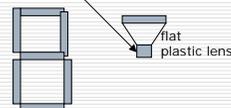
Light detector: excitation of valence electrons by photon absorption into the conduction band changes charge equilibrium

→ electric signal

CCD camera: two dimensional diode-array



LED: Light Emitting Diode



voltage brings extra charges into the conduction band – switch off leads to recombination → $E_{gap} = hf$

6. Major classes of materials

internal order (structure) and properties

- Gases

- Liquids

Mesomorphous systems:
Liquid crystals

- Solids (crystals)

6.1 Gases

ideal gas: no interaction – no bond formation – no „structure”

particle energy ~ kinetic energy

-> disorder?

two rules governing disorder

1. **average of kinetic energy** is determined by the temperature

$$\frac{1}{2} m \overline{v^2} = \frac{3}{2} kT$$

degree of freedom of motion for point-like particles

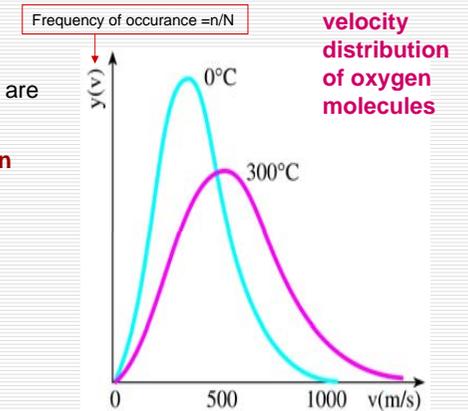
6.1 Gases

two rules

2. **deviations** from the average \overline{v} are determined by the

Maxwell – Boltzmann distribution function

Boltzmann distribution with $\mathcal{E} = 1/2 m v^2$



6.2 Liquids

particles ~ molecules built up by primary bonds

interaction between particles is by **secondary bonds**

	E_b (eV)
H-bond	several x 0.1
hydrophobic	~ 0.1
van der Waals dipole – point charge	~ 0.1-0.2
dipole – dipole	~ 0.02
dipole – induced dipole	~ 0.01
temporary dipoles	~ 0.02

$$\longleftrightarrow kT = 0.025 \text{ eV}$$

A large proportion of the secondary bonds are broken at room temperature

6.2 Liquids

interaction between particles is by **secondary bonds**

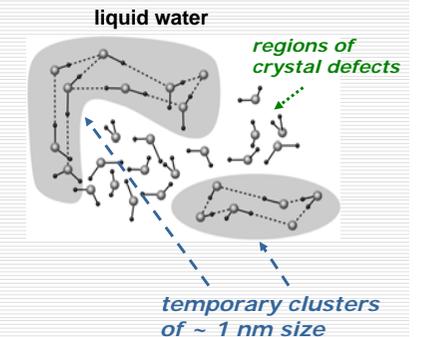
e.g. H-bonds in liquid water

$$\frac{n_{\text{broken}}}{n_0} \cong 1\% - 0.05\% \text{ at room tempr.}$$

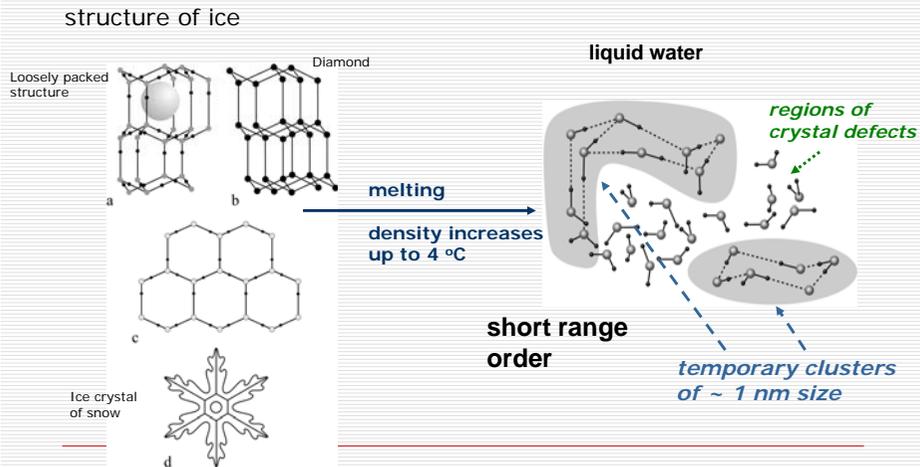
large number of broken bonds!

Properties:

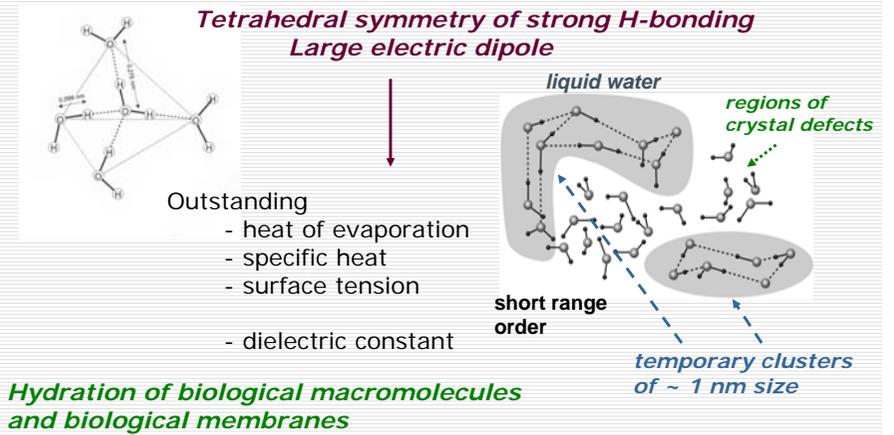
short range order
deformability
isotropy



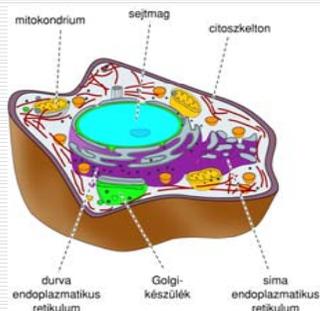
6.2 Liquids: the special properties of water



6.2 Liquids: the special properties of water



Live matter – what is its material class?



Hierarchy of strong and weaker bonds

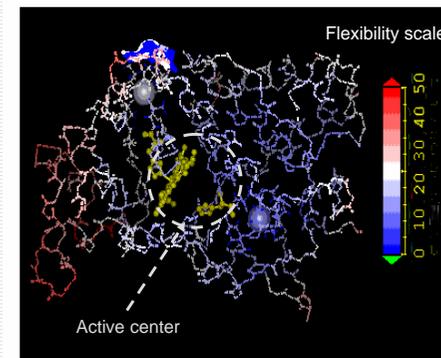
Primary bonds in base molecules, substrates, polymer chains – periodicity?? – quasicrystalline???,

Secondary bonds stabilize the secondary structure – sequential order, secondary order

Broad variety of bond strengths

Conformational dynamics: functional changes of the structure are made possible

Live matter ~ mesomorphous systems – intermediate between crystals and liquids

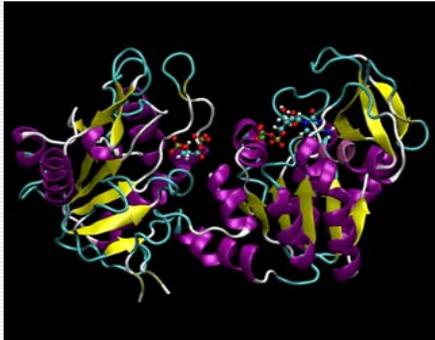


Hierarchy of strong and weaker bonds

Conformational dynamics make possible the opening of channels for substrates to reach the active center

Horseradish Peroxidase : polipeptide chain with two bound Ca ions and one prosthetic group (heme)

***Live matter ~ mesomorphous systems–
intermediate between crystals and liquids***



Phosphoglycerate kinase

Hierarchy of strong and weaker bonds

Conformational dynamics make possible the domain closure to bring ADP and 1,3-bisphosphoglycerate into contact

Mesomorphous systems – liquid crystals

Discussed next time – Dr. Voszka

Thank you for your attention
